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# Crystal Lattice Damage and Recovery of Rare-Earth Implanted Wide Bandgap Oxides

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#### ABSTRACT

Rare earth (RE) elements are important for the optical tuning of wide bandgap oxides (WBO) such as  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> or ZnO, because  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:RE or ZnO:RE show narrow emission lines in the visible, ultra-violet and infra-red region. Ion implantation is an attractive method to introduce dopant into the crystal lattice with an extraordinary control of the dopant ion composition and location, but it creates the lattice damage, which may render the dopant optically inactive. In this research work, we investigate the post-implantation crystal lattice damage of two matrices of wide-bandgap oxides,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and ZnO, implanted with rare-earth (RE) to a fluence of  $5 \times 10^{14}$ ,  $1 \times 10^{15}$  and  $3 \times 10^{15}$  atoms/cm<sup>2</sup>, and post-growth annealed in Ar and O<sub>2</sub> atmosphere, respectively. The effect of implantation and annealing on both crystal lattice damage caused by implantation with the same RE fluences in the case of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> seems to be higher than in the case of ZnO. Low temperature photoluminescence was used to investigate the optical activation of RE in both matrices after performed annealing.

**Keywords:** wide bandgap oxides, zinc oxide, gallium oxide, rare earth, ion implantation, Rutherford backscattering spectrometry, low temperature photoluminescence.

#### INTRODUCTION

Wide bandgap oxides (WBO) doped with rare earths (RE) have been extensively studied because of their possible application in phosphors, lasers, plasma displays, solar cells etc. [1]. Gallium oxide is a WBO of approximately 4.9 eV bandgap and high critical electrical field strength of 8 MV/cm. The wider bandgap of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> allows overcoming the thermal quenching effect, which is a problem in other semiconductors such as Si, ultimately enabling high temperature device operation, while large critical field allows highvoltage operation [2–5]. Due to the wide energy gap, photodetectors based on  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> are promising for solar-blind ultraviolet UV photodetectors [6]. Moreover, the high radiation hardness and chemical resistance of this material allow the devices to work even in harsh environment [7]. Zinc oxide, another WBO with a bandgap of 3.37 eV at T = 300 K [8], is a promising material for many optoelectronic and electronic applications such as solar cells, sensors, light emitting diodes and others. High quality single ZnO crystals produced on a large scale and the production of epitaxial films grown by efficient and cheaper industrial methods are additional advantages of this material. In case of luminescence,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> emits light in the UV and visible (blue, green) region [9–11], while ZnO emits light in the violet-blue region of the spectrum [12].

Both of these WBOs are advantageous for the photonic devices working in the visible and ultraviolet spectral region, however, the spectral emission region can be modified and enhanced by doping with a suitable dopant. Doping with RE elements can fulfill this purpose because of their characteristic sharp and narrow emission lines which may extend from infrared to ultraviolet region [13]. The origin of these emissions are the 4f intra-shell transitions, and high localization of the 4f electron shell helps to reduce the influence of the host material on RE luminescence [8]. Nevertheless, for RE in WBO, it is possible to obtain the effective resonant pumping of the f shell and thus enhancing emission spectra, which is important for applications in photonics. Because of this, both  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and ZnO are very interesting as host materials for RE elements.

Ion implantation is an attractive method for the introduction of dopant into the host material, because of the ability to introduce any ion into the target material with control of dopant concentration by using an appropriate dose and energy in the implantation process [14-12]. Moreover, a dopant can be introduced in to the matrix of a host material with concentration higher than its solid solubility limit. There are also indications of a higher photoluminescence (PL) intensity for a material implanted with RE as compared to the one in which RE was introduced into the host material during a growth process [15]. Unfortunately, ion implantation is a destructive technique, which causes crystal lattice damage. The process of irradiation and lattice damage will be even more complex in a geometrically complicated material, such as  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, with possibility of formation of many defect complexes [16]. Moreover, most of as-implanted RE ions are optically inactive [8,17].

The crystal lattice could be recovered by using different thermal treatments at the appropriate temperature and environment. For this purpose, various types of heat treatment have been used which could also optically activate RE [17, 18]. Rapid Thermal Annealing (RTA) is one of the most effective thermal treatments which leads to RE relocation and recovery of the crystal lattice [1]. Rutherford backscattering spectrometry in channeling mode (RBS/C) is the most suitable analytical technique for RE depth profiling and studying the lattice distortion caused by radiation damage. For the evaluation of the RBS/C spectra, the SIMNRA simulation code [19] is used to find out the RE concentration from RBS random spectra and its depth profile in the crystal lattice. The Monte Carlo CHAnneling SYmulation (McChasy) code [20, 21] is used to quantify the damage depth profile from the RBS spectra in the channeling mode. McChasy simulations code has the unique ability to simulate a mixture of defects: point defects modeled as randomly displaced atoms (RDA) and extended defects as dislocations, stacking faults, and dislocation loops (DIS), which contribute to the lattice distortion. Thus, as a result of the simulation, the depth distribution for both of these types of defects is evaluated separately.

In this paper, we compare lattice damage and recovery for two WBO materials,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and ZnO, implanted with Eu and Sm ions and dedicated to optoelectronic applications. The structural changes caused by ion implantation and optical activation of RE after thermal annealing are carefully investigated by RBS/C and low-temperature photoluminescence (LTPL).

## **EXPERIMENTAL DETAILS**

Commercial, single crystalline  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> substrates with the (010) orientation were provided by Kyma technologies, while ZnO epitaxial films were deposited on GaN/Al<sub>2</sub>O<sub>3</sub> substrates at 300°C by Atomic Layer Deposition (ALD). Further details of epitaxial ALD processes can be found in [1].

High-resolution X-ray diffraction (HRXRD) and RBS/C measurements were performed to investigate the crystalline quality of virgin samples before implantation. The good crystalline quality of both the materials was confirmed based on the sharp and intensive XRD signals from  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and from ZnO substrates. It should be noticed that in case of epitaxial ZnO layer, the signal was comparable to the signal coming from the GaN/Al<sub>2</sub>O<sub>3</sub> substrate. Reciprocal space maps (RSMs) measured in the vicinity of chosen reciprocal lattice spots for both ZnO and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> are shown in Figure 1 (a, b).

The  $\chi_{min}$  value, as calculated from the ratio of aligned to random backscattering yield of the RBS/C spectra, was evaluated as 1.4% for single  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal and as 5.3% for epitaxial ZnO/ GaN films used in this experiment. Implantation with Sm and Eu ions with energy of 150 keV to fluence of  $5 \times 10^{14}$ ,  $1 \times 10^{15}$  and  $3 \times 10^{15}$  atoms/cm<sup>2</sup> was performed at room temperature. The selected implantation energy allowed obtaining the 100 nm thick doped layer. This value was obtained from the McChasy simulations of the measured RBS



**Fig. 1.** Reciprocal Space Maps (RSM) of the symmetric (a) 00.2 XRD peak of ZnO, (b) 020 XRD peak of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>

spectra and is in agreement with the calculation using SRIM. Implantations were performed at the Institute of Physics, Maria Curie-Sklodowska University, Lublin, Poland, using an ion implanter equipped with arc discharge plasma ion [22]. RTA processing using an Accu Thermo AW-610 from Allwin21 Corporation system was performed at 800°C, in Ar for 0.5 min. for  $\beta$ -Ga<sub>2</sub>O<sub>2</sub>, and in O<sub>2</sub> for 10 min. in the case of ZnO. In this system the temperature of 800°C was achieved in 20 s, while the cooling down until room temperature was performed in a gas flow within about 300 s. The standard RBS/C measurements were performed at Helmholtz Zentrum Dresden-Rossendorf, Germany, with 1.7 MeV He- ions and using a Van de Graaff accelerator. In the RBS/C experiments, the silicon detector positioned at a scattering angle of  $170^{\circ}$  was used, with a depth resolution < 5 nmand an energy resolution < 20 keV. LTPL measurements of annealed ZnO:RE layers were performed at temperature T=5K with excitation energy  $E_{exc} = 3.49 \text{ eV}$ , while the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:RE samples were measured at T=10K with excitation energy  $E_{av} = 3.53$  eV. The McChasy simulations were performed based on experimental parameters obtained from the random RBS spectrum of the investigated sample as well as the aligned spectrum of the virgin sample. The detailed procedure of the McChasy simulations is described in [21].

#### RESULTS

#### **RBS/C** analysis

Figure 2a shows the random and <001> aligned RBS spectra of virgin ZnO and the same

sample implanted with different fluence of Sm ions. The low energy part of the spectra illustrates the signal of He<sup>2+</sup> ions backscattered from Zn, while the high energy part, above 1400 keV, demonstrates the signal of backscattered ions from RE. Implantation induced lattice damage can be assessed by the low energy part (leftpaneled Fig. 2a) and the location of RE in the matrix is determined from the RE part of spectra (right-paneled Fig. 2a). The solid lines show the simulated spectra calculated using McChasy computer code. Figure 2b shows a comparison of implanted ZnO:Sm samples before and after annealing. As can be seen in Figure 2a, the RE signal increases with increasing fluence and, correspondingly, lattice damage is build-up, which is visible as an increase in both dechanneling and the characteristic bulk damage peak observed in the aligned RBS spectra in the 1225–1310 keV energy region. After annealing, the damage near the surface almost disappeared, but the bulk damage peak is increased, which indicates that defects might have moved towards the bulk from the surface.

Figure 2c shows the random and <010> aligned RBS spectra of virgin and implanted  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> with different fluences of Sm. In this case, the low energy part of the spectra illustrates the He<sup>2+</sup> signal backscattered from Ga and the high energy part demonstrates the signal coming from RE. In the spectra of the implanted sample, a bimodal post-implanted damage profile is observed.

This kind of shape of damage peak observed in RBS/C spectra for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> implanted with RE has already been reported [13].



Fig. 2. Random (ran) and aligned (ali) RBS spectra for (a) ZnO implanted with Sm,
(b) ZnO implanted with Sm and annealed in O<sub>2</sub> at 800°C for 10 min., (c) β-Ga<sub>2</sub>O<sub>3</sub> implanted with Sm, and (d) β-Ga<sub>2</sub>O<sub>3</sub> implanted with Sm and annealed in Ar at 800°C for 30 sec. Solid lines in all the spectra show the McChasy simulations. The discrepancies between simulated lines vs experimental points, in Zn/Ga signal energy spectral range, are less than 5%

The damage peak starts to increase with increasing RE fluence, and dechanneling level is almost similar for all fluences used, however it is higher compared to the virgin sample. Fig. 2d shows the comparison of Sm- implanted and post-implantation annealed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> samples. Annealing leads to recovery of the lattice damage, which is evident from the decrease of the damage peak after annealing. Moreover, the defects appear to diffuse towards the surface, while the distribution of RE in the matrix remains unchanged. The calculated substitutional fraction,  $\mathbf{f}_{\mathbf{s}}$ , i.e. the relative amount of the impurity atoms at the lattice site positions, of RE as shown in Table 1 for  $\beta$ -Ga<sub>2</sub>O<sub>2</sub>:Sm is zero for all measured fluences, both before and after annealing, indicating that the RE atoms are located in the interstitial positions. In turn, for ZnO:RE, f clearly shows that impurity atoms are partially located in substitutional positions

after implantation, and they moved to interstitial positions only after annealing.

In order to obtain quantitative information about the damage, the McChasy simulations of the RBS/C spectra were performed. Figure 3 presents the depth distributions of RDA and DIS type of defects in ZnO:RE and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:RE obtained as a result of the best fit of the experimental RBS/C spectra shown in Figure 2 with the McChasy code.

The McChasy simulations reveal that in both ZnO:Sm (Fig. 3a) and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:Sm (Fig. 3b) the double damage peak occurs after implantation. For ZnO:Sm, with the fluence of 3 x 10<sup>15</sup> atoms/cm<sup>2</sup>, the atypical damage peak near the surface (called intermediate peak, IP) extends to ~25 nm, then a small depleted region and the typical bulk damage peak in the depth region from ~30nm to ~150 nm are observed. After annealing of ZnO:RE, as shown in Fig. 2b, the damage visible in the IP region has disappeared, while that observed in

Material	Nominal Conc. of RE (atoms/cm <sup>2</sup> )	Samples name	Real Conc. of RE (10 <sup>15</sup> atoms/cm <sup>2</sup> )	χ <sub>min</sub> (Ga)	χ <sub>min</sub> (Zn)	χ <sub>min</sub> (Sm)	f <sub>s</sub>
β-Ga₂O₃:Sm	0	Virgin	-	1.4	-	-	
	5e14	β-Ga₂O₃_Sm5e14 β-Ga₂O₃_Sm5e14_RTA	0.44	81.9 67.5	-	100 100	0 0
	1e15	β-Ga <sub>2</sub> O <sub>3</sub> _Sm1e15 β-Ga <sub>2</sub> O <sub>3</sub> _Sm1e15_RTA	0.96	89.6 78.6	-	100 100	0 0
	3e15	β-Ga₂O₃_Sm3e15 β-Ga₂O₃_Sm3e15_RTA	2.76	101.8 80.1	-	100 100	0 0
ZnO:Sm	Virgin	Virgin	-	-	5.3	-	-
	5e14	ZnO_Sm5e14 ZnO_Sm5e14_RTA	0.61 -	-	14.2 7.1	57 75	44 26
	1e15	ZnO_Sm1e15 ZnO_Sm1e15_RTA	1.00	-	14.7 11.5	61 100	40 0
	3e15	ZnO_Sm3e15 ZnO_Sm3e15_RTA	2.86 -	-	44.2 48.4	56 100	53 0

Table 1. Conc. of RE calculated by SIMNRA and substitutional fraction (f<sub>s</sub>) calculated by RBS/c spectra

the bulk region has increased. However, the Mc-Chasy simulations reveal that the total concentration of RDA before and after annealing is similar, which clearly shows that this type of defect near the surface is not lost but diffuses towards the bulk. DIS type of defects in ZnO:Sm remains unchanged before and after annealing.

The atypical damage peak near the surface shows ~100% of RDA which extends up to ~20 nm in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:Sm with the fluence of  $3 \times 10^{15}$ atoms/cm<sup>2</sup> (Fig. 3b). The depth damage, located close to the projected range of implanted ions, extends between 20 and 80 nm for RDA, while DIS type of defects are shifted up to ~110 nm. After annealing both concentration of RDA and DIS decrease illustrating the crystal lattice recovery. Moreover, in contrast to ZnO:RE system after annealing, the RDA defects diffuse to the surface in the case of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:Sm.

The ZnO:Eu systems (obtained with similar fluences and energy as for ZnO:Sm) were also studied. The defect distribution for ZnO:Sm and ZnO:Eu calculated from RBS/C spectra with the McChasy code exhibits similarity, as shown in Figure 3c. This result is consistent with our previous reports [23], according to which the damage level and profile, is identical for similar masses and fluences of implanted ions. Based on that, it could be assumed that defect distributions for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystals implanted with Eu and Sm should be similar.

#### Low-temperature photoluminescence

Typical low temperature PL spectra recorded in the visible region for both ZnO and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrices implanted with Eu and subsequently annealed is shown in Fig. 4(a, b). In the inset to Fig. 4a, the sharp near-band-edge emission (NBE) peak at 375 nm and a deep-level emission (DLE) centered around 510 nm, related to native defects, can be observed for virgin ZnO. This spectrum provides the background for the Eu peak that appears in the red spectral region, close to 600 nm, for the annealed ZnO:Eu sample.

The virgin  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystal shows a broad emission spectrum in the visible region (Fig. 4b), where the NBE is not visible. The annealed  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:RE crystal shows the Eu peak around 625 nm indicating the optical activation of RE after annealing. However, the signal of RE in  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is relatively weak, which means that the annealing conditions for this material need to be further optimized to enhance the Eu-related emission.

#### CONCLUSIONS

Rutherford Backscattering Spectrometry in the Channeling mode was used to assess the crystal lattice quality, lattice site location of RE, postimplantation damage and post-annealing recovery of the crystal structure of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and ZnO implanted with RE. The SIMNRA and McChasy simulations of experimental RBS/C spectra reveal that the implantation process is very different for ZnO and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrices.

The implanted RE ions are situated in the interstitial positions in the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrix and stay the same after annealing. In the case of ZnO, the substitutional fraction varies between

40 and 53% depending on the fluence and decreases after annealing. The McChasy simulations confirm different defect behavior in both



Fig. 3. Depth distributions of (a) RDA+DIS defects for implanted and annealed ZnO:Sm
(b) RDA for implanted and annealed β-Ga<sub>2</sub>O<sub>3</sub> :Sm + ZnO:Sm (c) RDA+DIS defects for ZnO:Eu and ZnO:Sm obtained by McChasy Simulations

matrices after annealing, which also suggests different mechanisms of defect creation and diffusion. For implanted to the fluence of  $3 \times 10^{15}$ atoms/cm<sup>2</sup>  $\beta$ -Ga<sub>2</sub>O<sub>2</sub>:RE, the annealing in the Ar atmosphere at 800°C for 30 s removes defects in the implanted region, while for ZnO implanted with RE with the same fluence (and similar dpa) annealed in O<sub>2</sub> at 800 °C for 10 mins the concentration of defect remains unchanged, the defect depth profile is only changing. According to our unpublished studies, the defects distribution in ZnO:RE is almost independent of the annealing conditions. However, it should be pointed out that the annealing atmospheres strongly influence the luminescence efficiency. In this aspect, β-Ga<sub>2</sub>O<sub>2</sub>:RE crystals require further investigation, because the observed luminescence response, although visible, is not very high. It is expected that the luminescence efficiency should be higher for WBO with the wider band gap, so



**Fig. 4**. LT PL spectra for (a) ZnO and ZnO:Eu annealed in  $O_2$  at 800 °C for 10 min and (b)  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:Eu annealed in Ar at 800 °C for 30 sec

our future studies will be focused on finding the optimal annealing conditions for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>:RE.

What is interesting, contrary to widespread belief, the presented results show that the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> matrix is less radiation resistant compared to the ZnO matrix. For  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> implanted with RE the amorphization level is reached for the fluence of  $3 \times 10^{15}$  atoms/cm<sup>2</sup> (20 dpa), while for ZnO the amorphization level cannot be achieved even after ions bombardment with the fluence up to 100 dpa. [24]

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